A MODIFIED QUASI-REVERSIBILITY METHOD FOR A CLASS OF ILL-POSED CAUCHY PROBLEMS

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Abstract. The goal of this paper is to present some extensions of the method of quasi-reversibility applied to an ill-posed Cauchy problem associated with an unbounded linear operator in a Hilbert space. The key point to our proof is the use of a new perturbation to construct a family of regularizing operators for the considered problem. We show the convergence of this method, and we estimate the convergence rate under a priori regularity assumptions on the problem data.

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1. Introduction and Motivation

Throughout this paper H denotes a complex Hilbert space endowed with the inner product $(.,.)$ and the norm $\|.\|, \mathcal{L}(H)$ stands for the Banach algebra of bounded linear operators on H.

Let A be a linear unbounded operator with dense domain $\mathcal{D}(A)$. Assume that A is self-adjoint, positive definite in H . We consider the abstract final value problem

$$
u'(t) + Au(t) = 0, \ \ 0 < t < T, \ \ u(T) = f,\tag{FVP}
$$

where f is a given function in H .

Such problems are not well-posed in the Hadamard sense [15]; the solution (if it exists) does not depend continuously on the data.

Physically, problems of this nature arise in different contexts. Beyond their interest in connection with standard diffusion problems (then A is usually the Laplace operator $-\Delta$), they also appear, for instance, in some deconvolution problems, such as deblurring processes $[6](A$ is often a fractional power of $-\Delta$), in hydrology [16, 36] and also in many other practical applications of mathematical physics and engineering sciences.

Since the semigroup $S(t) = e^{-tA}$ is not time-inverted, to obtain a well-posed problem, we would like to find an operator $S_{\alpha}(t)$, $\alpha > 0$, "close to" $S(t)$ in some sense for which the final value problem (FVP) is well-posed.

In the mathematical literature various methods have been proposed for solving backward Cauchy problems. We can notably mention the method of quasisolution of Tikhonov $[37]$, the method of quasi-reversibility of Lattes and Lions [23], the method of *logarithmic convexity* $[1, 7, 21, 25, 30]$, the *iterative proce*dures of Kozlov and Maz'ya $[5, 22]$, the quasi boundary value method $[8, 11, 20]$, the C-regularized semigroups theory [3, 26, 27, 33].

In the method of quasi-reversibility, the main idea consists in replacing A in (FVP) by $A_{\alpha} = g_{\alpha}(A)$. In the original method [23] Lattès and Lions proposed $g_{\alpha}(A) = A - \alpha A^2$, and in the Gajewski and Zaccharias quasi-reversibility method [12, 14, 19, 34], $g_{\alpha}(A) = A(I + \alpha A)^{-1}$, to obtain a well-posed problem in the backward direction. Then, using the information from the solution of the perturbed problem, another well-posed problem is constructed, and its solution sometimes can be taken to be an approximate solution of the original ill-posed problem (FVP) .

Difficulties may arise when using the method of quasi-reversibility discussed above. an essential difficulty is that the operator coefficient is replaced by an operator of second order, which makes the numerical implementation rather difficult; in addition, the error $(e(\alpha))$ introduced by a small changes in the final value f is of order $e^{\frac{1}{\alpha}}$. For these reasons, we propose a modified quasireversibility based on the perturbation

$$
A_{\alpha} = g_{\alpha}(A) = -\frac{1}{pT} \log \left(\alpha + e^{-pT A} \right), \quad \alpha > 0, \quad p \ge 1.
$$

An advantage of this new perturbation lies in the fact that it is bounded ($A_{\alpha} \in$ $\mathcal{L}(H)$, which gives the well-posedness in the forward and backward direction for the perturbed problem, while another advantage is that this perturbation provides the best possible approximate solution, while the amplification factor of the error resulting from the approximated problem is better as compared with other results. We note that our approach generalizes many results obtained by other methods.

2. Preliminaries and Basic Results

In this section, we give the notation and functional is needed in the sequel.

If $B \in \mathcal{L}(H)$, we denote by $\mathcal{N}(B)$ the kernel of B and by $\mathcal{R}(B)$ the range of B. We denote by $\{E_\lambda, \lambda \geq \gamma > 0\}$ the spectral resolution of the identity associated to A, and by $S(t) = e^{-tA}$ ∞ γ $e^{-t\lambda} dE_{\lambda} \in \mathcal{L}(H), t \geq 0$, the C_0 semigroup generated by $-A$. Some basic properties of $S(t)$ are listed in the following theorem:

Theorem 2.1 (see [32], Ch. 2, Theorem 6.13, p. 74). For the family of operators $S(t)$, the following properties are valid:

- (1) $||S(t)|| \leq 1, \quad \forall t \geq 0;$
- (2) the function $t \mapsto S(t)$, $t > 0$, is analytic;
- (3) for every real $r \geq 0$ and $t > 0$, the operator $S(t) \in \mathcal{L}(H, \mathcal{D}(A^{r}))$;
- (4) for every integer $k \ge 0$ and $t > 0$, $||S^{(k)}(t)|| = ||A^k S(t)|| \le c(k)t^{-k}$;
- (5) for every $x \in \mathcal{D}(A^r)$, $r \geq 0$ we have $S(t)A^r x = A^r S(t)x$.

Theorem 2.2. For every $t > 0$, the operator $S(t)$ is self-adjoint and one-toone with dense range $(S(t) = S(t)^*, \mathcal{N}(S(t)) = \{0\}$ and $\overline{\mathcal{R}(S(t))} = H)$.

Proof. A is self-adjoint and since $S(t)^* = (e^{-tA})^* = e^{-tA^*} = e^{-tA}$, we have $S(t)^* = S(t)$. Let $h \in \mathcal{N}(S(t_0))$, $t_0 > 0$, then $S(t_0)h = 0$, which implies that $S(t)S(t_0)h = S(t + t_0)h = 0$, $t \ge 0$. Using analyticity, we obtain that $S(t)h = 0, t \geq 0$. Strong continuity at 0 now gives $h = 0$. This shows that $\mathcal{N}(S(t_0)) = \{0\}.$ By

$$
\overline{\mathcal{R}(S(t_0))} = \mathcal{N}(S(t_0))^{\perp} = \{0\}^{\perp} = H
$$

we conclude that $\mathcal{R}(S(t_0))$ is dense in H. \Box

For more details, we refer the reader to a general version of Theorem 2.2 in the case of analytic semigroups in Banach spaces (Lemma 2.2, [9]).

Remark 2.1. If we replace A by $B = pA$ in Theorem 2.2, we obtain $\mathcal{N}(S(pt)) =$ ${0}$ and $\overline{\mathcal{R}(S(pt))} = H, p > 0, t > 0.$

Remark 2.2 (Smoothing effect and irreversibility). By Theorems 2.1 and 2.2, we observe that the solution of the direct Cauchy problem:

$$
u'(t) + Au(t) = 0, \ \ 0 < t \leq T, \quad u(0) = u_0,
$$

has the following smoothing effect: admitting the initial value $u(0)$ to belong only to H, for all $t > 0$ we obtain

$$
\mathcal{R}(S(t)) \subset \mathcal{C}^{\infty}(A) \stackrel{def}{=} \cap_{n=1}^{\infty} \mathcal{D}(A^n)
$$

(a space more regular than H , see [13]). It follows that for the final value problem (FVP) to have a solution, we should have $u(T) \in C(A) \subseteq \mathcal{R}(S(T)),$ where $\mathcal{C}(A)$ is an admissible class for which the FVP is solvable. This shows that this problem is irreversible in the sense

$$
S(T-t): H \to \mathcal{R}(S(T-t)) \subset \mathcal{C}^{\infty}(A) \subsetneq H, \quad 0 \le t < T,
$$

and $\mathcal{R}(S(T-t)) \neq \overline{\mathcal{R}(S(T-t))}$; in other words, $S(T-t)^{-1} = S(t-T) \notin \mathcal{L}(H)$.

For notational convenience and simplicity, we set

$$
\mathcal{C}_{\theta}(A) = \{ \phi \in H : \ \|\phi\|_{\theta}^{2} = \int_{\gamma}^{+\infty} e^{2T\theta\lambda} d\|E_{\lambda}\phi\|^{2} < +\infty \}, \quad \theta \ge 0.
$$

Evidently,

$$
\mathcal{C}_{\theta_2}(A) \subseteq \mathcal{C}_{\theta_1}(A), \quad \theta_2 \ge \theta_1.
$$

We also need the following technical lemma.

Lemma 2.1. For $x \geq 0$ and $\tau \in [0,1]$, we have

$$
(1+x)^{\tau} - 1 \leq \tau x (1+x)^{\tau} (1+\tau x)^{-1}.
$$

The proof is justified by a simple differential calculus (see [24], p. 101.)

Remark 2.3. The operational calculus for self-adjoint operators and Lemma 2.1 plays the key role in our analysis and calculations.

3. The Approximate Problem

In this section, we give a constructive method based on the quasi-reversibility approach to construct a stable approximate solution to the ill-posed problem $(FVP).$

Description of the method:

Step 1. Let v_{α} be a solution of the following perturbed problem

$$
v'_{\alpha}(t) + A_{\alpha}v_{\alpha}(t) = 0, \quad 0 \le t < T, \quad v_{\alpha}(T) = f,
$$
 (FVP) _{α}

where the operator A is replaced by A_{α} .

Step 2. We use the initial value

$$
v_{\alpha}(0) = \varphi_{\alpha}
$$

in the problem

$$
u'_{\alpha}(t) + Au_{\alpha}(t) = 0, \ \ 0 < t \leq T, \quad u_{\alpha}(0) = \varphi_{\alpha}.
$$
 (IVP)_{\alpha}

Step 3. Finally, we show that

$$
\Phi_{\alpha}(f) = ||u_{\alpha}(T) - f|| \longrightarrow 0 \text{ as } \alpha \longrightarrow 0.
$$

4. Analysis of the Method and Main Results

4.1. Analysis of A_{α} and its consequences. We begin our study by giving some qualitative properties of A_{α} .

For $0 < \alpha \leq \alpha_* = 1 - e^{-\gamma T}$, $p \geq 1$, we define

$$
A_{\alpha} = g_{\alpha}(A) := -\frac{1}{pT} \log \left(\alpha + e^{-pT A} \right) = \int_{\gamma}^{\infty} -\frac{1}{pT} \log \left(\alpha + e^{-pT \lambda} \right) dE_{\lambda}.
$$

For more details concerning the logarithm of operators and its spectral properties see, e.g., [4, 17, 29].

Proposition 4.1. We have

(1)
$$
A_{\alpha} \in \mathcal{L}(H)
$$
 and $||A_{\alpha}|| \le \frac{1}{pT} \log \left(\frac{1}{\alpha}\right)$;
\n(2) $A_{\alpha} = A_{\alpha}^{*} \ge 0$ and $A_{\alpha}A^{\theta}v = A^{\theta}A_{\alpha}v$, $v \in \mathcal{D}(A^{\theta})$, $\theta \ge 0$;
\n(3) $\forall v \in \mathcal{D}(A)$, $\lim_{\alpha \to 0} ||A_{\alpha}v - Av|| = 0$;
\n(4) $\forall v \in H$, $S_{\alpha}(t)v = e^{-tA_{\alpha}}v \longrightarrow S(t)v$ as $\alpha \longrightarrow 0$.

Proof. 1. The boundedness of the operator A_{α} follows immediately from the properties of $g_{\alpha}(\lambda)$, $\lambda \geq \gamma$. Indeed, by the choice of α we have

$$
\alpha + e^{-pT\gamma} \le \alpha + e^{-T\gamma} \le 1,
$$

which implies that

$$
\underline{g_{\alpha}}=g_{\alpha}(\gamma)=-\frac{1}{pT}\log\left(\alpha+e^{-pT\gamma}\right)\geq0
$$

and

$$
\overline{g_{\alpha}} = \lim_{\lambda \to +\infty} g_{\alpha}(\lambda) = -\frac{1}{pT} \log(\alpha) > 0.
$$

Observing that $g'_{\alpha}(\lambda) > 0$, we have $g_{\alpha} \nearrow$ and $\sup g_{\alpha}(\lambda) = \overline{g_{\alpha}}$. By using a λ≥γ spectral representation of A_{α} and the preceding remarks, we obtain the desired result.

2. Via the H-functional calculus (see [10]) and the self-adjointness of A we establish (2).

3. Let $v \in \mathcal{D}(A)$, we have

$$
A_{\alpha}v = -\frac{1}{pT}\log\left(\alpha + e^{-pTA}\right)A^{-1}Av.
$$

Let us denote

$$
B_{\alpha} = -\frac{1}{pT} \log \left(\alpha + e^{-pTA} \right) A^{-1} = \int_{\gamma}^{+\infty} M_{\alpha}(\lambda) dE_{\lambda},
$$

where $M_{\alpha}(\lambda) = -\frac{1}{n^2}$ $\frac{1}{pT} \log \left(\alpha + e^{-pT\lambda} \right) \lambda^{-1}$. By the definition of α , we observe that $M_{\alpha}(\lambda) \geq 0$ for all $\lambda \geq \gamma$ and

$$
M_{\alpha}(\lambda) = 1 - \frac{1}{pT} \log \left(1 + \alpha e^{pT\lambda} \right) \lambda^{-1} \ge 0
$$

which implies that $M_{\alpha}(\lambda) \leq 1$ for all $\lambda \geq \gamma$. Consequently, the operator B_{α} is uniformly bounded, i.e., $||B_{\alpha}|| \leq 1$, \forall 0 < $\alpha \leq 1 - e^{-\gamma T}$.

Let $v = e^{-pT A} h$, $h \in H$, we have

$$
||(B_{\alpha} - I)v||^2 = \int_{\gamma}^{+\infty} \left(\frac{1}{pT} \log(1 + \alpha e^{pT\lambda})\lambda^{-1}\right)^2 e^{-2pT\lambda} d||E_{\lambda}h||^2.
$$
 (a)

Since $log(1 + x) \leq x$, $x \geq 0$, then the quantity (*a*) can be estimated as follows

$$
||(B_{\alpha}-I)v||^2 \leq \left(\frac{\alpha}{p\gamma T}\right)^2 ||h||^2 \longrightarrow 0, \text{ as } \alpha \longrightarrow 0,
$$

from which it follows that $B_\alpha v \longrightarrow v$ in H as $\alpha \longrightarrow 0$, $\forall v \in \mathcal{R}(S(pT))$. But $\mathcal{R}(S(pT))$ is dense in H and B_{α} is uniformly bounded on H, hence, by continuity,

 $\forall v \in H, B_{\alpha}v \longrightarrow v, \text{ as } \alpha \longrightarrow 0.$

In particular, for $v \in \mathcal{D}(A)$ we obtain

$$
A_{\alpha}v = B_{\alpha}Av \longrightarrow Av, \text{ as } \alpha \longrightarrow 0.
$$

4. Since $A_{\alpha} \in \mathcal{L}(H)$, we can define

$$
S_{\alpha}(t) = e^{-tA_{\alpha}} = \left(\alpha + e^{-pTA}\right)^{\frac{t}{pT}} = \sum_{n=0}^{+\infty} \frac{(-t)^n}{n!} A_{\alpha}^n, \quad t \in \mathbb{R}.
$$

It is obvious to see that $||S_\alpha(t)|| \leq 1$, $t \geq 0$. Thus $S_\alpha(t)$, $t \geq 0$ is a strongly continuous semigroup of contraction on H. Also, $\frac{d}{dt}$ $\frac{d}{dt}S_{\alpha}(t) = -A_{\alpha}(t)S_{\alpha}(t)$ and

$$
S_{\alpha}(t) - S_{\beta}(t) = \int_{0}^{t} \frac{d}{d\tau} \Big(S_{\beta}(t-\tau) S_{\alpha}(\tau) \Big) d\tau
$$

=
$$
\int_{0}^{t} \Big(S_{\beta}(t-\tau) S_{\alpha}(\tau) (A_{\beta} - A_{\alpha}) \Big) d\tau.
$$

Then $\forall t \geq 0, 0 < \alpha, \beta \leq 1 - e^{-\gamma T}, h \in \mathcal{D}(A)$, we have

$$
||S_{\alpha}(t)h - S_{\beta}(t)h|| \leq t||A_{\beta}h - A_{\alpha}h||,
$$

which shows that $\{S_{\alpha}(t)h\}$ is a Cauchy sequence in H, uniformly in $t \in [0, T]$ (by virtue of (3) in Proposition 4.1).

To complete the proof of (4), observe that $S_{\alpha}(t)$ is a contraction and $\mathcal{D}(A)$ is dense in H , so the limit

$$
S_{\alpha}(t)h \longrightarrow \widetilde{S}(t)h \text{ as } \alpha \longrightarrow 0, \quad t \ge 0,
$$

extends to all $h \in H$, and holds uniformly in $t \in [0, T]$. It is clear that $\widetilde{S}(t) \in$ $\mathcal{L}(H)$ is a strongly continuous semigroup of contraction on H.

Let $h \in \mathcal{D}(A), t > 0$, then

$$
||S(t)h - S_{\alpha}(t)h|| = \left\| \int_{0}^{t} \frac{d}{d\tau} \left(S(\tau)S_{\alpha}(t-\tau)h \right) d\tau \right\|
$$

$$
\leq \int_{0}^{t} ||S_{\alpha}(t-\tau)(A-A_{\alpha})S(\tau)h|| d\tau
$$

$$
\leq \int_{0}^{t} ||(A-A_{\alpha})S(\tau)h|| d\tau.
$$

Now we use

$$
||A_{\alpha}S(\tau)h|| = ||B_{\alpha}AS(\tau)h|| \le ||S(\tau)Ah||
$$

to get

$$
||(A_{\alpha} - A)S(\tau)h|| \leq 2||S(\tau)Ah||.
$$

Since $||S(\tau)Ah||$ is continuous, by the dominated convergence theorem we have

$$
\lim_{\alpha \to 0} ||S(t)h - S_{\alpha}(t)h|| \le \int_{0}^{t} \lim_{\alpha \to 0} ||(A - A_{\alpha})S(\tau)h|| d\tau = 0,
$$

which implies that $S_{\alpha}(t) \longrightarrow S(t) = \widetilde{S}(t)$ on $\mathcal{D}(A)$ as $\alpha \longrightarrow 0$. According to the density of $\mathcal{D}(A)$ in H we conclude that $S_{\alpha}(t) \longrightarrow S(t) = \widetilde{S}(t)$ on H as $\alpha \longrightarrow 0$.

We note here that by a direct calculation with the help of Lemma 2.1 we can show

$$
\forall h \in H, \quad S_{\alpha}(t)h \longrightarrow S(t)h \text{ as } \alpha \longrightarrow 0.
$$

Indeed, let $v = e^{-ptA}h$, $h \in H$, we have

$$
||S_{\alpha}(t)v - S(t)v||^{2} = \int_{\gamma}^{+\infty} \left(\left(\alpha + e^{-pT\lambda} \right)^{\frac{t}{pT}} - e^{-t\lambda} \right)^{2} e^{-2pT\lambda} d||E_{\lambda}h||^{2}.
$$

By virtue of Lemma 2.1 and $\alpha + e^{-pT\lambda} \leq 1$, the function

$$
M_{\alpha}(\lambda) = \left(\alpha + e^{-pT\lambda}\right)^{\frac{t}{pT}} - e^{-t\lambda} = e^{-t\lambda} \left(\left(1 + \alpha e^{pT\lambda}\right)^{\frac{t}{pT}} - 1\right)
$$

can be estimated as follows:

$$
M_{\alpha}(\lambda) \le \frac{\frac{t}{pT} \alpha e^{pT\lambda} \left(\alpha + e^{-pT\lambda}\right)^{\frac{t}{pT}}}{\left(1 + \frac{t}{pT} \alpha e^{pT\lambda}\right)} \le \frac{\alpha}{p} e^{pT\lambda}.
$$

From this we derive

$$
||S_{\alpha}(t)v - S(t)v||^{2} = \int_{\gamma}^{+\infty} M_{\alpha}(\lambda)^{2} e^{-2pT\lambda} d||E_{\lambda}h||^{2} \le \left(\frac{\alpha}{p}\right)^{2} ||h||^{2} \longrightarrow 0 \text{ as } \alpha \longrightarrow 0.
$$

According to the density of $\mathcal{R}(S(pT))$ in H and $||S_{\alpha}(t)|| \leq 1, t \geq 0$, we conclude that

$$
\forall h \in H, \quad S_{\alpha}(t)h \longrightarrow S(t)h \text{ as } \alpha \longrightarrow 0.
$$

The proof of Proposition 4.1 is complete. \Box

4.2. Convergence results. Now we are ready to state and prove the main results of this paper.

It is useful to know exactly the admissible set for which the problem (FVP) has a solution. The following lemma answers this question.

Lemma 4.1 (see [8], Lemma 1). The problem (FVP) has a solution if and only if $f \in C_1(A)$, and its unique solution is represented by

$$
u(t) = e^{(T-t)A}f.
$$
\n⁽¹⁾

By using the semi-groups theory and the properties of $S_{\alpha}(t)$ we have the following theorems.

Theorem 4.1. For all $f \in H$, the function

$$
v_{\alpha} = e^{(T-t)A_{\alpha}} f = \left(\alpha + e^{-pTA}\right)^{-\frac{(T-t)}{pT}} f \tag{2}
$$

is a unique solution of the problem $(FVP)_{\alpha}$ and it depends continuously on f.

To show the continuous dependence of v_{α} on f, we note that

$$
||v_{\alpha}(t)|| = \left\| \left(\alpha + e^{-pTA} \right)^{-\frac{(T-t)}{pT}} f \right\|
$$

$$
\leq \left(\frac{1}{\alpha} \right)^{\frac{T-t}{pT}} ||f|| \leq \left(\frac{1}{\alpha} \right)^{\frac{1}{p}} ||f|| = e(\alpha) ||f||. \tag{3}
$$

From (2) we construct

$$
\varphi_{\alpha} = v_{\alpha}(0) = \left(\alpha + e^{-pTA}\right)^{-\frac{1}{p}} f.
$$

Theorem 4.2. The problem $(IVP)_{\alpha}$ is well-posed, and its solution is represented by

$$
u_{\alpha}(t) = S(t)\varphi_{\alpha} = e^{-tA} \left(\alpha + e^{-pTA}\right)^{-\frac{1}{p}}f.
$$
 (4)

Theorem 4.3. For all $f \in H \|u_\alpha(T) - f\| \longrightarrow 0$ as $\alpha \longrightarrow 0$.

Proof. We compute

$$
||u_{\alpha}(T) - f||^2 = \int\limits_{\gamma}^{+\infty} H_{\alpha}(\lambda)^2 d||E_{\lambda}f||^2,
$$
\n(5)

where

$$
H_{\alpha}(\lambda) = \frac{\left((\alpha + e^{-pT\lambda})^{\frac{1}{p}} - e^{-T\lambda}\right)}{\left(\alpha + e^{-pT\lambda}\right)^{\frac{1}{p}}} = \frac{e^{-T\lambda}\left((\alpha e^{pT\lambda} + 1)^{\frac{1}{p}} - 1\right)}{\left(\alpha + e^{-pT\lambda}\right)^{\frac{1}{p}}}.
$$

If we put $x = \alpha e^{pT\lambda}, \tau = \frac{1}{n}$ $\frac{1}{p}$, then by virtue of Lemma 2.1, the function $H_{\alpha}(\lambda)$ can be estimated as α

$$
H_{\alpha}(\lambda) \le \frac{\alpha}{\alpha + pe^{-pT\lambda}}.\tag{6}
$$

From (6) it follows

$$
||u_{\alpha}(T) - f||^2 \le \int_{\gamma}^{+\infty} \left\{ \frac{\alpha}{\alpha + pe^{-pT\lambda}} \right\}^2 d||E_{\lambda}f||^2.
$$
 (7)

Fix $\varepsilon > 0$, and choose N so that $+\infty$ N $d||E_\lambda f||^2 < \frac{\varepsilon}{2}$ $\frac{\varepsilon}{2}$. Thus

$$
||u_{\alpha}(T) - f||^2 \leq \int_{\gamma}^{N} \left\{ \frac{\alpha}{\alpha + pe^{-pT\lambda}} \right\}^2 d||E_{\lambda}f||^2
$$

+
$$
\int_{N}^{+\infty} \left\{ \frac{\alpha}{\alpha + pe^{-pT\lambda}} \right\}^2 d||E_{\lambda}f||^2,
$$
 (8)

which gives

$$
|u_{\alpha}(T) - f||^2 \le \left(\frac{\alpha}{p}\right)^2 e^{2pTN} ||f||^2 + \frac{\varepsilon}{2}.
$$
\n(9)

So, by taking α such that $\left(\frac{\alpha}{\alpha}\right)$ p e^{2pTN} $\|f\|^2 < \frac{\varepsilon}{2}$ $\frac{\varepsilon}{2}$, we end the proof. \Box

Remark 4.1. Here we do not have a convergence rate.

Theorem 4.4. If $f \in C_{\theta}(A)$, $p \geq 1$, $0 < \theta < 1$, then $||u_{\alpha}(T) - f||$ converges to zero with order $\alpha^{\frac{\theta}{p}}$.

Proof. We compute

$$
||u_{\alpha}(T) - f||^2 = \int_{\gamma}^{+\infty} \left\{ \frac{H_{\alpha}(\lambda)}{e^{\theta T \lambda}} \right\}^2 e^{2\theta T \lambda} d||E_{\lambda} f||^2 \le \int_{\gamma}^{+\infty} G_{\alpha}(\lambda)^2 e^{2\theta T \lambda} d||E_{\lambda} f||^2, (10)
$$

where

$$
G_{\alpha}(\lambda) = \frac{\alpha}{\left(\alpha + pe^{-pT\lambda}\right)e^{\theta T\lambda}} > 0.
$$

Differentiation with respect to λ yields

$$
G'_{\alpha}(\lambda) = \frac{\alpha T e^{\theta T \lambda} \left(p(p - \theta) e^{-pT\lambda} - \theta \alpha \right)}{\left((\alpha + p e^{-pT\lambda}) e^{\theta T \lambda} \right)^2}.
$$

Thus $G'_{\alpha}(\lambda) = 0$ if $\lambda = \lambda^* = \frac{1}{\gamma^2}$ $p\bar{T}$ $\log \Big($ $p(p - \theta)$ $\left(\frac{\partial}{\partial \alpha} - \theta\right)$. Since $G'_{\alpha}(\lambda) > 0$ if $\lambda < \lambda^*$, $G'_{\alpha}(\lambda) < 0$ if $\lambda > \lambda^*$ and $\lim_{\lambda \to +\infty} G_{\alpha}(\lambda) = 0$, we have that λ^* is the critical value at which G_{α} achieves its maximum. Thus we have the inequality

$$
G_{\alpha}(\lambda) \le G_{\alpha}(\lambda^*) = c(p, \theta) \alpha^{\frac{\theta}{p}}, \qquad (11)
$$

where $c(p, \theta) = \left(\begin{array}{c} 1 \end{array} \right)$ 1 p $\int_{\frac{p+\theta}{p}}^{\frac{p+\theta}{p}} (p-\theta)^{\frac{p-\theta}{p}} \theta^{\frac{\theta}{p}} \leq 1.$

Combining (10) and (11), we arrive at

$$
||u_{\alpha}(T) - f|| \le c(p, \theta) \alpha^{\frac{\theta}{p}} ||f||_{\theta}.
$$
\n(12)

Noting that in the case $1 \leq p \leq \theta$, we have the estimate

$$
||u_{\alpha}(T) - f|| \leq \alpha ||f||_{\theta}.
$$
\n(13)

Theorem 4.5. For all $f \in H$, the problem (FVP) has a solution u if and only if the sequence $\varphi_{\alpha} = u_{\alpha}(0)$ converges in H. Furthermore, we then have $u_{\alpha}(t)$ converges to $u(t)$, as α tends to zero uniformly in t.

Proof. Assume that
$$
\lim_{\alpha \to 0} \varphi_{\alpha} = \varphi_0
$$
 exists. Let $w(t) = S(t)\varphi_0$. We compute
\n
$$
||w(t) - u_{\alpha}(t)|| = ||S(t)\varphi_0 - S(t)\varphi_{\alpha}|| = ||S(t)(\varphi_0 - \varphi_{\alpha})||
$$
\n
$$
\le ||\varphi_0 - \varphi_{\alpha}||.
$$

This implies

$$
\sup_{0\leq t\leq T}||w(t)-u_{\alpha}(t)||\leq ||\varphi_0-\varphi_{\alpha}||\longrightarrow 0, \text{ as } \alpha\longrightarrow 0.
$$

Since $\lim_{\alpha \to 0} u_{\alpha}(T) = f$ and $\lim_{\alpha \to 0} u_{\alpha}(T) = w(T)$, we infer that $w(T) = f$. Thus, $w(t) = S(t)\varphi_0$ solves the problem (FVP) and satisfies the condition $w(T) = f$. Now let us assume that $u(t)$ is a solution of (FVP) . From Lemma 4.1 we

have $u(0) = S(-T)f \in H$, i.e., $||u(0)||^2 = ||f||_1^2 =$ $+\infty$ γ $e^{2T\lambda} d||E_{\lambda}f||^2 < \infty$. Let $N > 0$ and $\varepsilon > 0$ such that $+\infty$ N $e^{2T\lambda} d||E_\lambda f||^2 < \frac{\varepsilon}{2}$ 2 . We compute

$$
||u_{\alpha}(0) - u(0)||^{2} = \int_{\gamma}^{+\infty} F_{\alpha}^{2}(\lambda) d||E_{\lambda}f||^{2},
$$

where

$$
F_{\alpha}(\lambda) = e^{T\lambda} - \left(\alpha + e^{-pT\lambda}\right)^{-\frac{1}{p}}.
$$

By simple calculations with the help of Lemma 2.1, $F_{\alpha}(\lambda)$ can be estimated as

$$
F_{\alpha}(\lambda) \leq \left(\frac{\alpha e^{pT\lambda}}{p + \alpha e^{pT\lambda}}\right) e^{T\lambda} = K_{\alpha}(\lambda) e^{T\lambda}.
$$

Then

$$
||u_{\alpha}(0) - u(0)||^{2} \le \int_{\gamma}^{+\infty} K_{\alpha}^{2}(\lambda) e^{2T\lambda} d||E_{\lambda}f||^{2} \le I_{1} + I_{2},
$$

where

$$
I_1 = \int_{\gamma}^{N} K_{\alpha}^2(\lambda) e^{2T\lambda} d\|E_{\lambda}f\|^2 \le \left(\frac{\alpha}{p}\right)^2 e^{2(p+1)TN} \|f\|^2,
$$

$$
I_2 = \int_{N}^{+\infty} K_{\alpha}^2(\lambda) e^{2T\lambda} d\|E_{\lambda}f\|^2 < \frac{\varepsilon}{2}.
$$

Now if we choose α such that $\left(\frac{\alpha}{n}\right)$ p $\sqrt{2}$ $e^{2(p+1)TN}$ || f ||² < $\frac{\varepsilon}{2}$ $\frac{\varepsilon}{2}$, then we have $||u_{\alpha}(0) - u(0)||^{2} < \varepsilon.$

This shows that

$$
u_{\alpha}(0) \longrightarrow u(0)
$$
 as $\alpha \longrightarrow 0$.

Theorem 4.6. If $f \in C_{1+\theta}(A)$, $p \ge 1$, $0 < \theta < 1$, then $||u_{\alpha}(0) - u(0)||$ converges to zero with order $\alpha^{\frac{\theta}{p}}$.

Proof. By similar calculations to those used in Theorems 4.4 and 4.5, we have

$$
||u_{\alpha}(0) - u(0)||^2 \le \int_{\gamma}^{+\infty} K_{\alpha}(\lambda)^2 e^{-2\theta T \lambda} e^{2(1+\theta)T\lambda} d||E_{\lambda}f||^2
$$

$$
\le \int_{\gamma}^{+\infty} G_{\alpha}(\lambda)^2 e^{2(1+\theta)T\lambda} d||E_{\lambda}f||^2
$$

$$
\le (G_{\alpha}^{\infty})^2 ||f||_{1+\theta}^2,
$$

with $G_{\alpha}^{\infty} = \sup_{\lambda \geq \gamma}$ $G_{\alpha}(\lambda) \leq c(p,\theta)\alpha^{\frac{\theta}{p}}$ (see estimate (11)).

From Theorems 4.5–4.6 we obtain

Corollary 4.1. If $f \in C_{1+\theta}(A)$, $\theta > 0$, then $||u_{\alpha}(t) - u(t)||$ converges to zero with order $\alpha^{\frac{\theta}{p}}$ uniformly in t.

We end this paper by constructing a family of regularizing operators for $(FVP).$

Definition 4.1. A family $\{R_{\alpha}(t), \alpha > 0, t \in [0, T]\} \subset \mathcal{L}(H)$ is called a family of regularizing operators for the problem (FVP) if for each solution $u(t)$, $0 \leq t \leq T$ of the (FVP) with final element f, and for any $\delta > 0$, there exists $\alpha(\delta) > 0$, such that

$$
\alpha(\delta) \longrightarrow 0, \quad \delta \longrightarrow 0, \tag{R_1}
$$

$$
||R_{\alpha(\delta)}(t)f_{\delta} - u(t)|| \longrightarrow 0, \quad \delta \longrightarrow 0,
$$
 (R₂)

for each $t \in [0, T]$ provided that f_δ satisfies $||f_\delta - f|| \leq \delta$.

Define $R_{\alpha}(t) = e^{-tA} \left(\alpha + e^{-pTA} \right)^{-\frac{1}{p}}, t \geq 0, \alpha > 0$; it is clear that $R_{\alpha}(t) \in$ $\mathcal{L}(H)$.

In the following we will show that $R_{\alpha}(t)$ is a family of regularizing operators for the (FVP) .

Theorem 4.7. Assuming that $f \in C_1(A)$, then (\mathcal{R}_2) holds.

Proof. We have

$$
H_{\alpha}(t) = ||R_{\alpha}(t)f_{\delta} - u(t)|| \le ||R_{\alpha}(t)(f_{\delta} - f)|| + ||R_{\alpha}(t)f - u(t)||
$$

= $\Delta_1(t) + \Delta_2(t),$ (14)

where

$$
\Delta_1(t) = ||R_\alpha(t)(f_\delta - f)|| \le \left(\frac{1}{\alpha}\right)^{\frac{1}{p}} \delta,\tag{15}
$$

and

$$
\Delta_2(t) = ||R_\alpha(t)f - u(t)||. \tag{16}
$$

Choose $\alpha =$ √ δ, then $\alpha(\delta)$ → 0, δ → 0, and

$$
\Delta_1(t) \le \delta^{\frac{2p-1}{2p}} \longrightarrow 0 \text{ as } \delta \longrightarrow 0. \tag{17}
$$

Now, by virtue of Theorem 4.5 we have

$$
\Delta_2(t) = \|u_\alpha(t) - u(t)\| \longrightarrow 0 \text{ as } \delta \longrightarrow 0,
$$
\n(18)

uniformly in t . Combining (17) and (18) we obtain

$$
\sup_{0 \le t \le T} \|R_{\alpha}(t)f_{\delta} - f\| \longrightarrow 0 \text{ as } \delta \longrightarrow 0. \tag{19}
$$

This shows that $R_{\alpha}(t)$ is a family of regularizing operators for the (FVP) . \Box

We proceed with the same technique used to establish the preceding results. We show that

$$
\Delta_2(t) \le C(p, t, T) \alpha^{\frac{t}{pT}} \|f\|_1, \quad t > 0,
$$
\n(20)

with

$$
C(p,t,T) = p^{-\frac{pT+t}{pT}} (pT-t)^{\frac{pT-t}{pT}} t^{\frac{t}{pT}} T^{-1} \le 1.
$$

Example 4.1. We give an example to clarify our study. As an example we consider the following backward heat equation:

$$
\frac{\partial u(x,t)}{\partial t} - \frac{\partial^2 u(x,t)}{\partial x^2} = 0, \quad x \in (0,\pi), \quad t \in (0,T),
$$

\n
$$
u(0,t) = u(\pi,t) = 0, \quad t \in (0,T),
$$

\n
$$
u(x,T) = \varphi(x), \qquad x \in [0,\pi].
$$
\n(BHE)

Here $u(x, t)$ represents the temperature at time t at a point x of a thin metal wire of length π .

The problem (BHE) can be formulated in the abstract form as follows

$$
u'(t) + Au(t) = 0, \quad 0 < t < T, \quad u(T) = \varphi,
$$

where the linear operator

$$
A: D(A) \subset H = L_2(0, \pi) \longrightarrow L_2(0, \pi)
$$

is defined by

$$
A:=-\frac{\partial^2}{\partial x^2}
$$

with

$$
\mathcal{D}(A) := H^2(0, \pi) \cap H_0^1(0, \pi).
$$

It is easy to show that the operator A is self-adjoint and positive with discrete spectrum $(\sigma(A) = {\lambda_n}_{n\geq 1})$. We denote by $\{e_n\}_{n\geq 1}$ the orthonormal eigenbasis in H, associated to the eigenvalues $\{\lambda_n\}_{n\geq 1}$ such that:

$$
Ae_n = \lambda_n e_n, \quad \lambda_n = n^2, \quad e_n(x) = \sqrt{\frac{2}{\pi}} \sin(nx), \quad n = 1, 2, \cdots
$$

$$
(e_n, e_m) = \int_0^{\pi} e_n(x) e_m(x) dx = \frac{2}{\pi} \int_0^{\pi} \sin(nx) \sin(mx) dx = \begin{cases} 1, & n = m, \\ 0, & n \neq m. \end{cases}
$$

$$
\forall f \in L_2(0, \pi), \quad f(x) = \sum_{n \ge 1} c_n(f) e_n(x) = \sum_{n \ge 1} \frac{2}{\pi} \int_0^{\pi} f(y) \sin(ny) dy \sin(nx).
$$

By the functional calculus for the self-adjoint operator we have

$$
g(A) = \int_{\sigma(A)} g(\lambda) dE_{\lambda} = \sum_{n \ge 1} g(\lambda_n) P_{\lambda_n}, \quad P_{\lambda_n}(\xi) = c_n(\xi) e_n,
$$

for each Borel-measurable function $g : \sigma(A) \longrightarrow \mathbb{C}$.

In particular, for g_{α} we obtain the perturbation

$$
A_{\alpha} = g_{\alpha}(A) = \sum_{n \ge 1} -\frac{1}{pT} \log \left(\alpha + e^{-pT\lambda_n} \right) P_{\lambda_n}
$$

and the associated semigroup

$$
S_{\alpha}(t) = e^{-tA_{\alpha}} = \sum_{n \ge 1} \left(\alpha + e^{-pT\lambda_n} \right)^{\frac{t}{pT}} P_{\lambda_n}
$$

.

In this case, the approximate solution $u_{\alpha}(x, t)$ takes the form

$$
u_{\alpha}(x,t) = R_{\alpha}(t)f(x) = \frac{2}{\pi} \sum_{n \ge 1} e^{-tn^2} \left(\alpha + e^{-pTn^2}\right)^{\frac{-1}{p}} \int_{0}^{\pi} f(y) \sin(ny) dy \sin(nx).
$$

Concluding remarks and generalization.

1. From (3) we observe that the error factor $e(\alpha)$ introduced by small changes in the final value f is of order $(\frac{1}{2})$ α $\big)^{\frac{1}{p}}, p \geq 1.$

2. In the case $p = 1$ the representation of u_{α} . coincides with that obtained in the method developed in [8], with the same error factor $e(\alpha)$ of order $\frac{1}{\alpha}$.

α **3.** In [11] (resp. [14, 23]) the error factor $e(\alpha)$ is of order $\frac{1}{\alpha(1+\log(T/\alpha))}$ (resp. $e^{\frac{1}{\alpha}}$).

Observing that for $p > 1$, $\left(\frac{1}{p}\right)$ α $\sqrt{1}$ $\frac{1}{p} < e^{\frac{1}{\alpha}},$ (1) α $\sqrt{1}$ $\frac{1}{p} < \frac{1}{\alpha(1 + \log(T/\alpha))}$ for $\alpha > 0$. This shows that our approach has a nice regularizing effect and gives a better

approximation as compared with the methods developed in [8, 11, 14, 23].

Let us consider

$$
u_t(t) + Au(t) = 0, \quad 0 < t < T, \quad u(0) = f,\tag{CP}_1
$$

$$
u_t(t) + Au(t) = 0, \quad 0 < t < T, \quad u(T) = f,
$$
 (CP)₂

where A is a self-adjoint, linear unbounded operator in H and changes the sign with $0 \in \rho(A)$ $(A^{-1}$ exists and $A^{-1} \in \mathcal{L}(H)$).

The spectral theory of self-adjoint operators enables us to write \mathcal{L}

$$
h = \int_{\mathbb{R}} dE_{\lambda} h = \int_{\mathbb{R}_{-}} dE_{\lambda} h + \int_{\mathbb{R}_{+}} dE_{\lambda} h = h_{-} + h_{+}, \quad h \in H,
$$

i.e., the Hilbert space H decomposes into the direct sum $H = H_-\oplus H_+$, and α α

$$
A = \int_{\mathbb{R}} \lambda \, dE_{\lambda} = \int_{\mathbb{R}_{-}} \lambda \, dE_{\lambda} + \int_{\mathbb{R}_{+}} \lambda \, dE_{\lambda} = A_{-} + A_{+}.
$$

It is well known that the problem $(CP)_1$ (resp. $(CP)_2$) is incorrectly posed in the sense of Hadamard. In order to regularize this problem, we propose the following family of operators:

$$
R_{\alpha}(t) = e^{(T-t)A_{-}} \left(\alpha e^{pTA_{-}} + (1 - \alpha) \right)^{-\frac{1}{p}}
$$

$$
+ e^{-tA_{+}} \left(\alpha + (1 - \alpha) e^{-pTA_{+}} \right)^{-\frac{1}{p}}
$$

$$
= R_{\alpha}^{-}(t) + R_{\alpha}^{+}(t), \quad 0 < \alpha < 1, \quad p \ge 1.
$$

We conclude our study with the following lemma (fundamental lemma).

Lemma 4.2. If $f = f_+ + f_+ \in H$, then $(CP)_1$ (resp. $(CP)_2$) has a solution if and only if \overline{y}

$$
\int_{\mathbb{R}_-} e^{-2T\lambda} d||E_{\lambda}f_-||^2 < +\infty \ \ (resp. \ \int_{\mathbb{R}_+} e^{2T\lambda} d||E_{\lambda}f_+||^2 < +\infty).
$$

Proof. If $\int e^{-2T\lambda} d||E_{\lambda}f_{-}||^{2} < +\infty$, we define $u(t) = e^{-tA_{-}}f_{-} + e^{-tA_{+}}f_{+}$. It is R−

not difficult to verify that $u(t)$ satisfies $(CP)_1$. Conversely, let $u(t)$ be a solution to $(CP)_1$. Then $u(T) = h = h_{-} + h_{+} = e^{-TA_{-}}f_{-} + e^{-TA_{+}}f_{+} \in H$. This implies that $h = e^{-TA} - f = H$, i.e., $||h||^2 =$ ϵ R− $e^{-2T\lambda} d||E_{\lambda}f_{-}||^{2} < +\infty$. In a similar

way we show the second part of the lemma . \Box

We define $u_{\alpha}(t) = R_{\alpha}(t) f$. With the same methodology used to establish the results of convergence in the case where A is positive, we show that

$$
u_{\alpha}(T) \longrightarrow f \text{ as } \alpha \longrightarrow 0,
$$
\n(21)

$$
u_{\alpha}(0) \longrightarrow f \text{ as } \alpha \longrightarrow 1,
$$
\n(22)

and $R_{\alpha}(t)$ is a family of regularizing operators for $(CP)_1$ (resp. $(CP)_2$).

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